

Rayleigh-Taylor instability in confined porous media: pore-scale simulations and experiments

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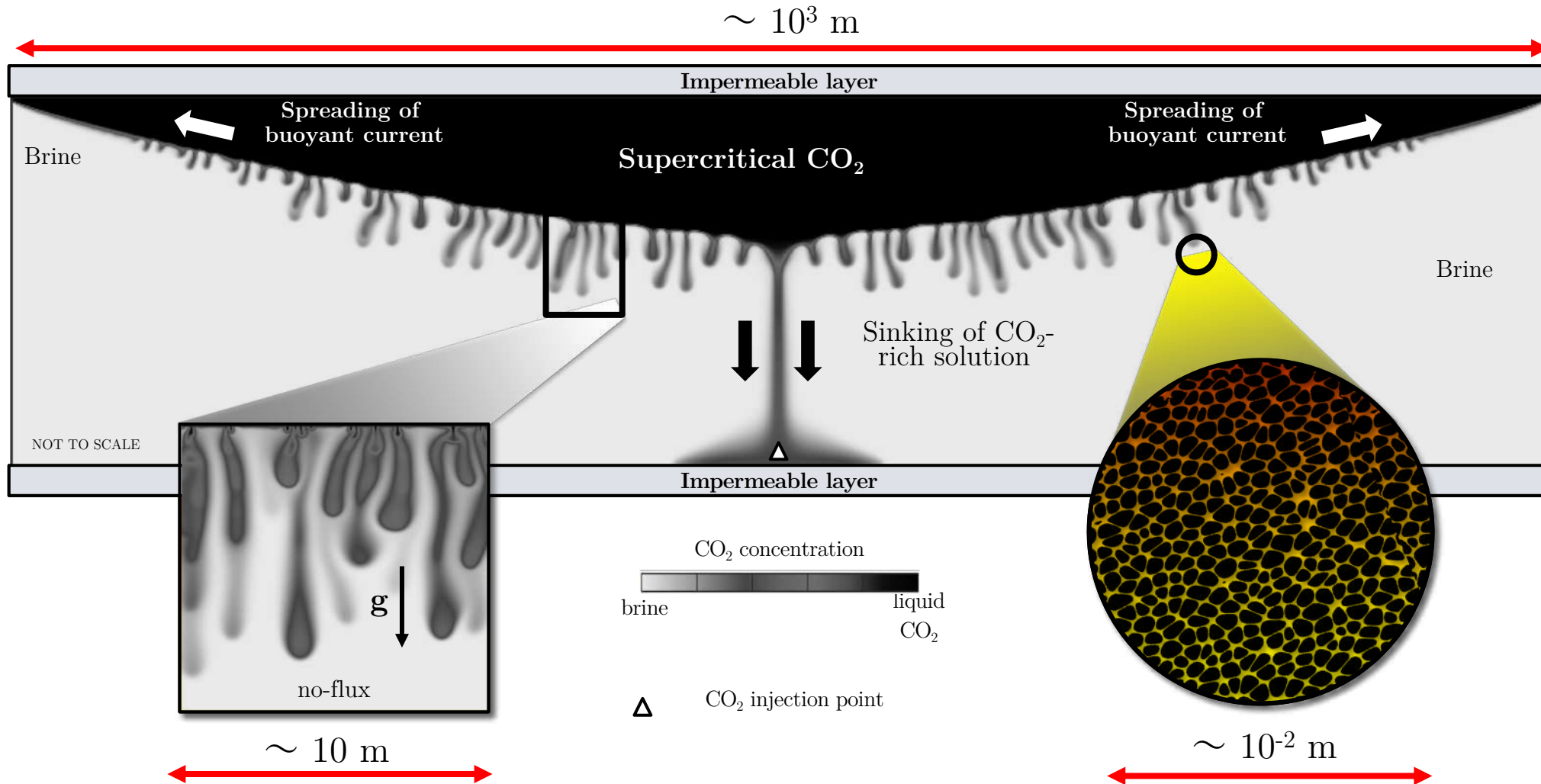
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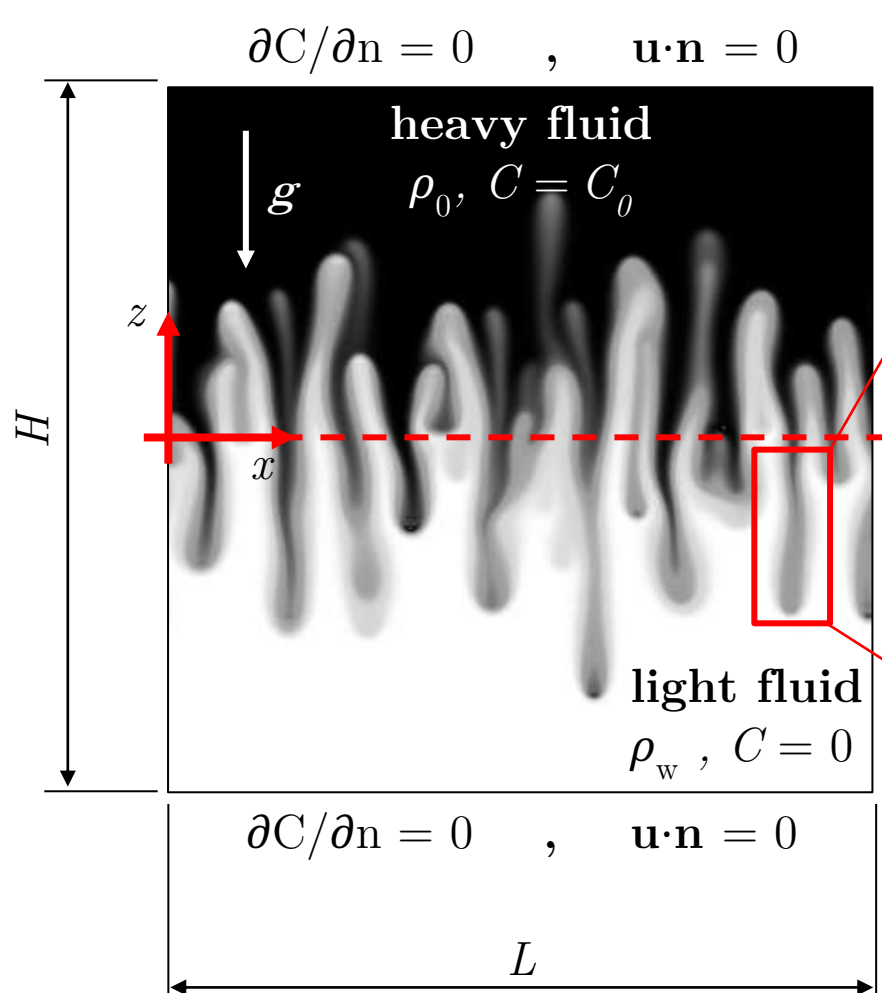
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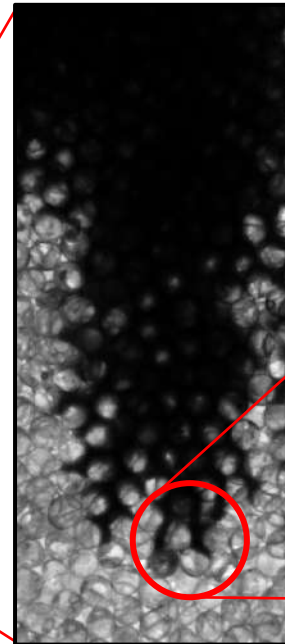
De Paoli, *Phys. Fluids*. (2021)

Rayleigh-Taylor instability in confined porous media

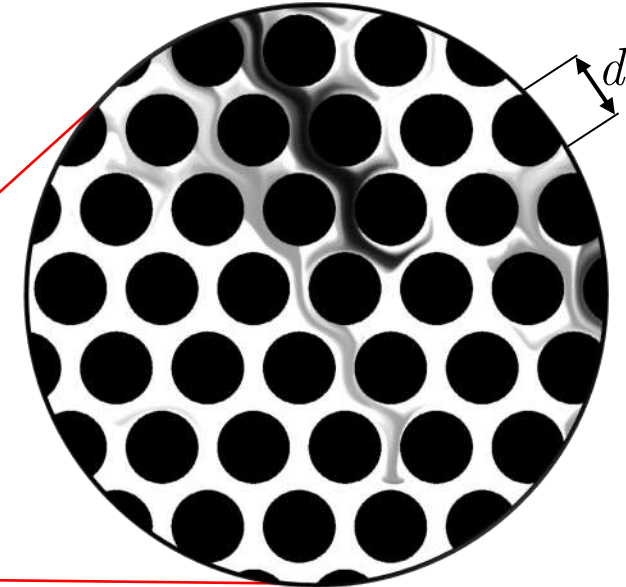
Marco De Paoli, Physics of Fluids Group, University of Twente



experiments

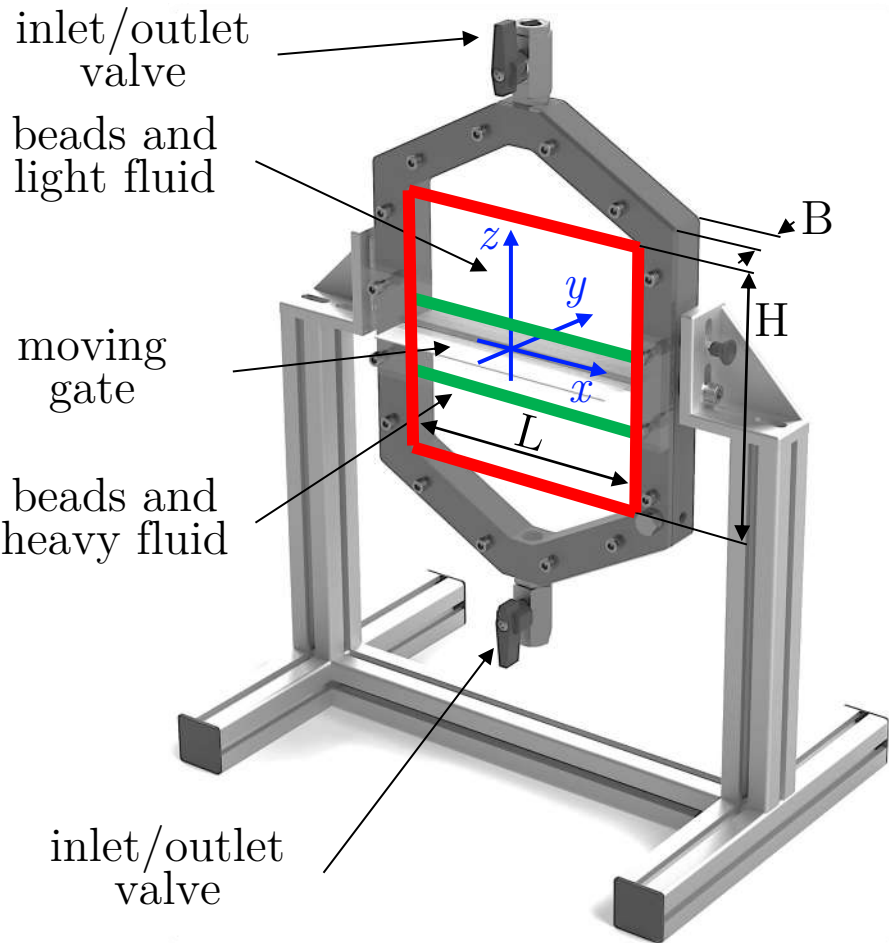


simulations

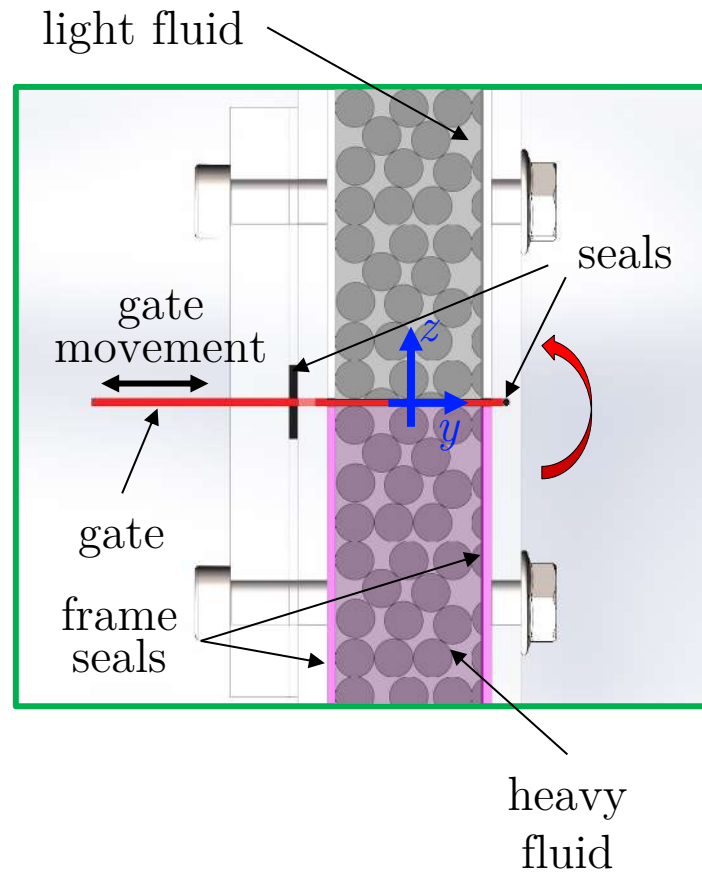


- High Schmidt number $Sc = \frac{\mu}{\rho_0 D}$
- Porosity matched $\phi = 0.37$
- Solid impermeable to solute
- Linear dependency $\rho(C)$

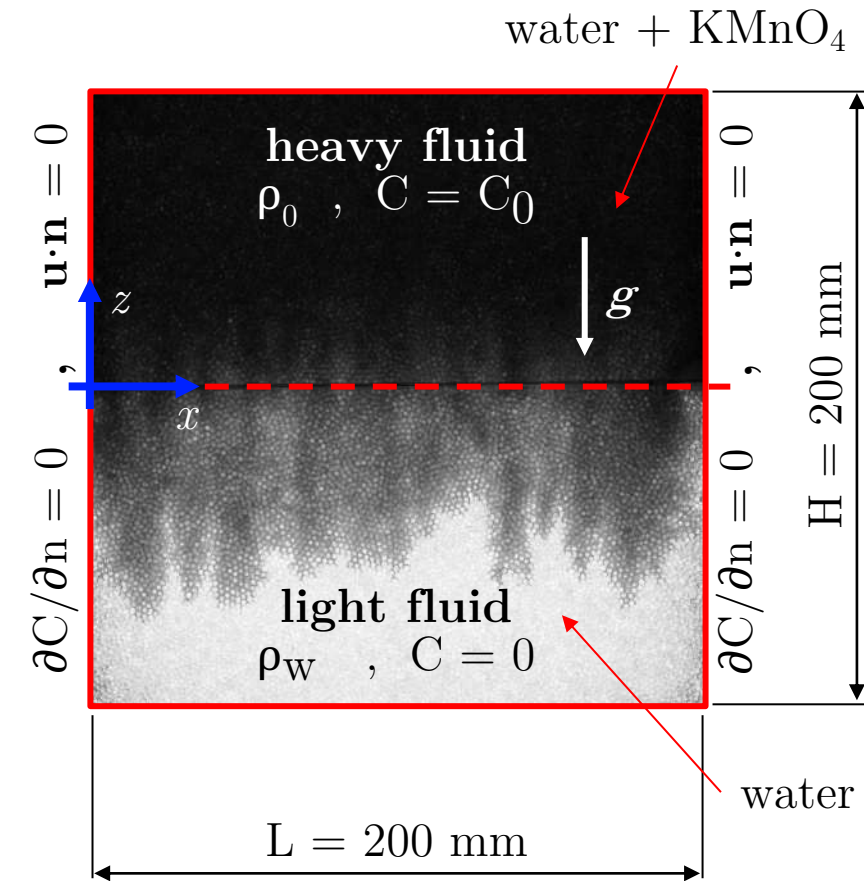
(a) Hele-Shaw cell

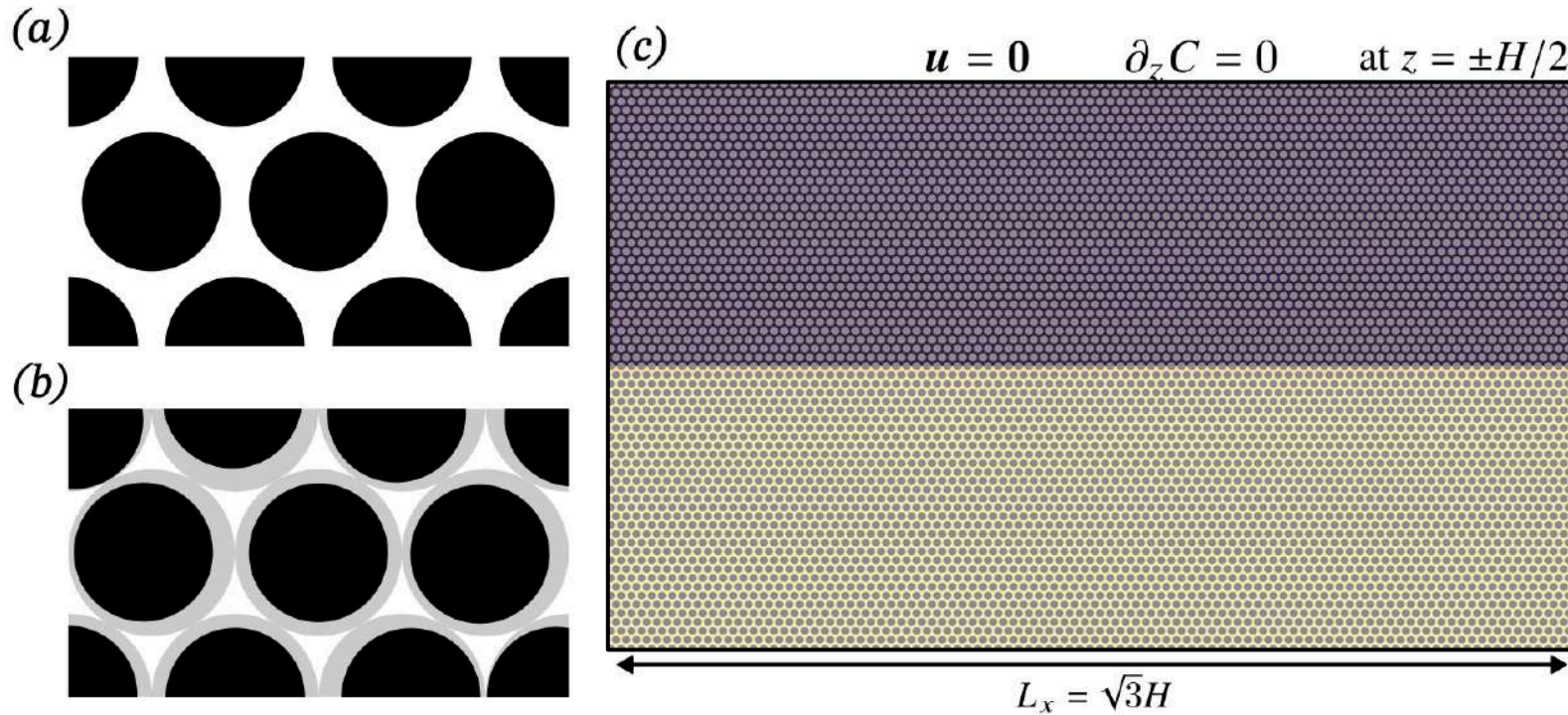


(b) gate (side view)



(c) measurement region



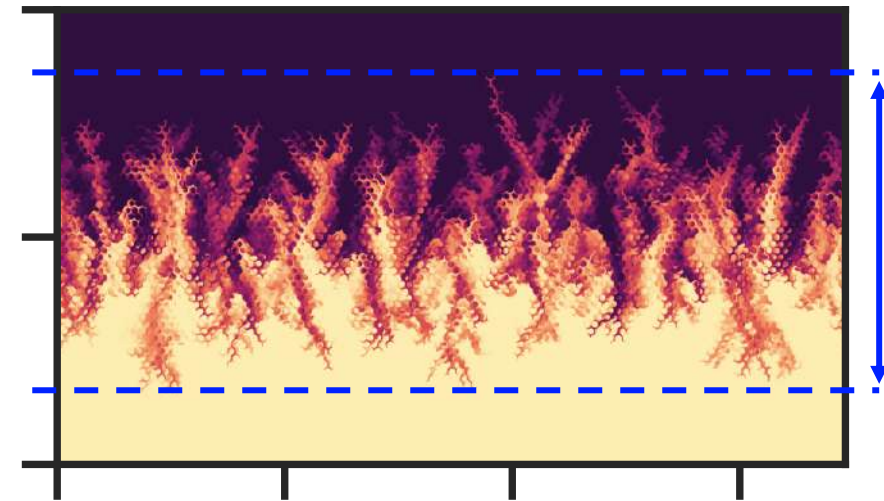


Finite difference
(AFiD, open source)
+
Immersed Boundaries
Method

Resolution:

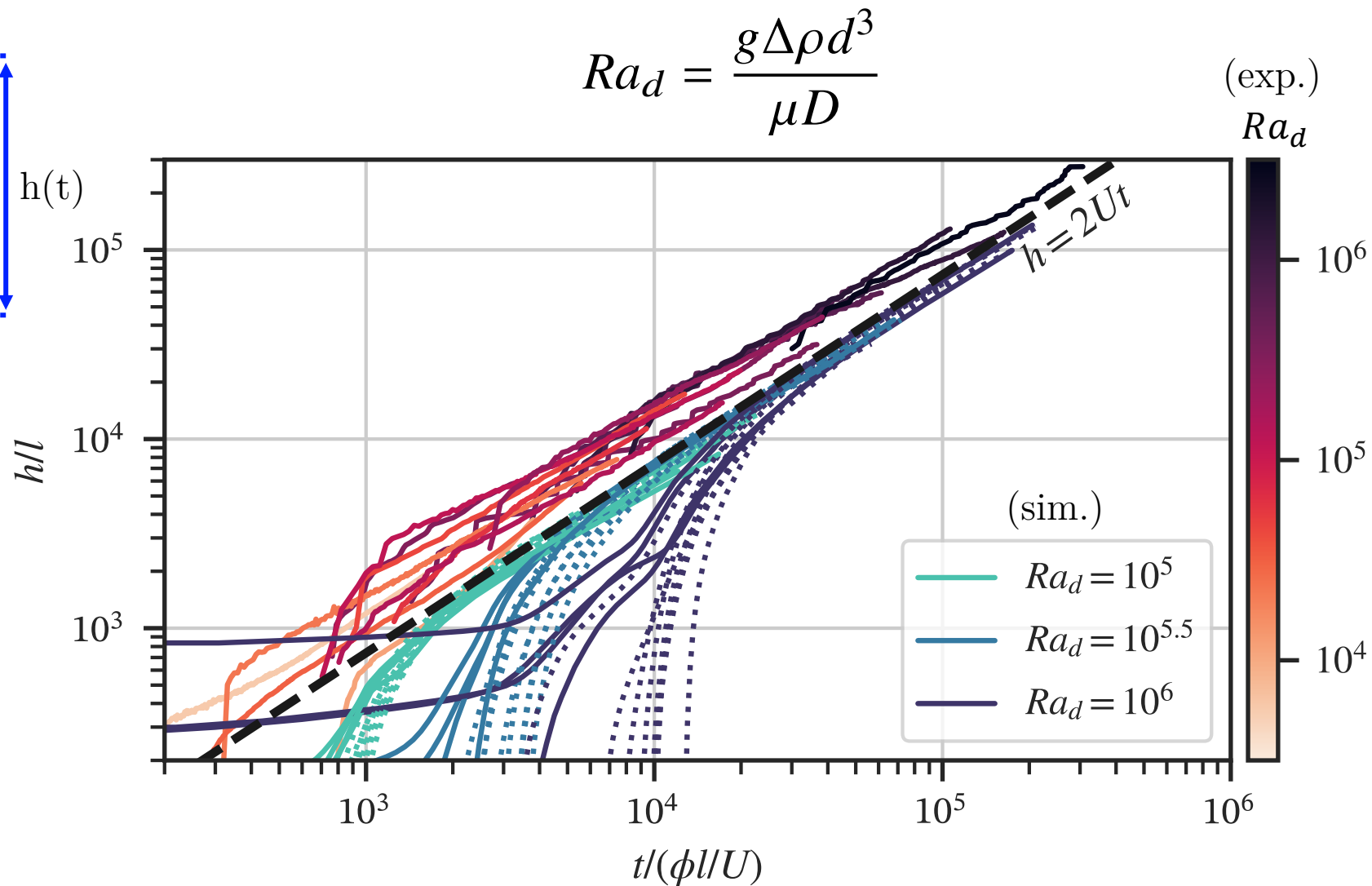
- velocity: ≥ 32 points per diameter
- conc.: ≥ 128 points per diameter

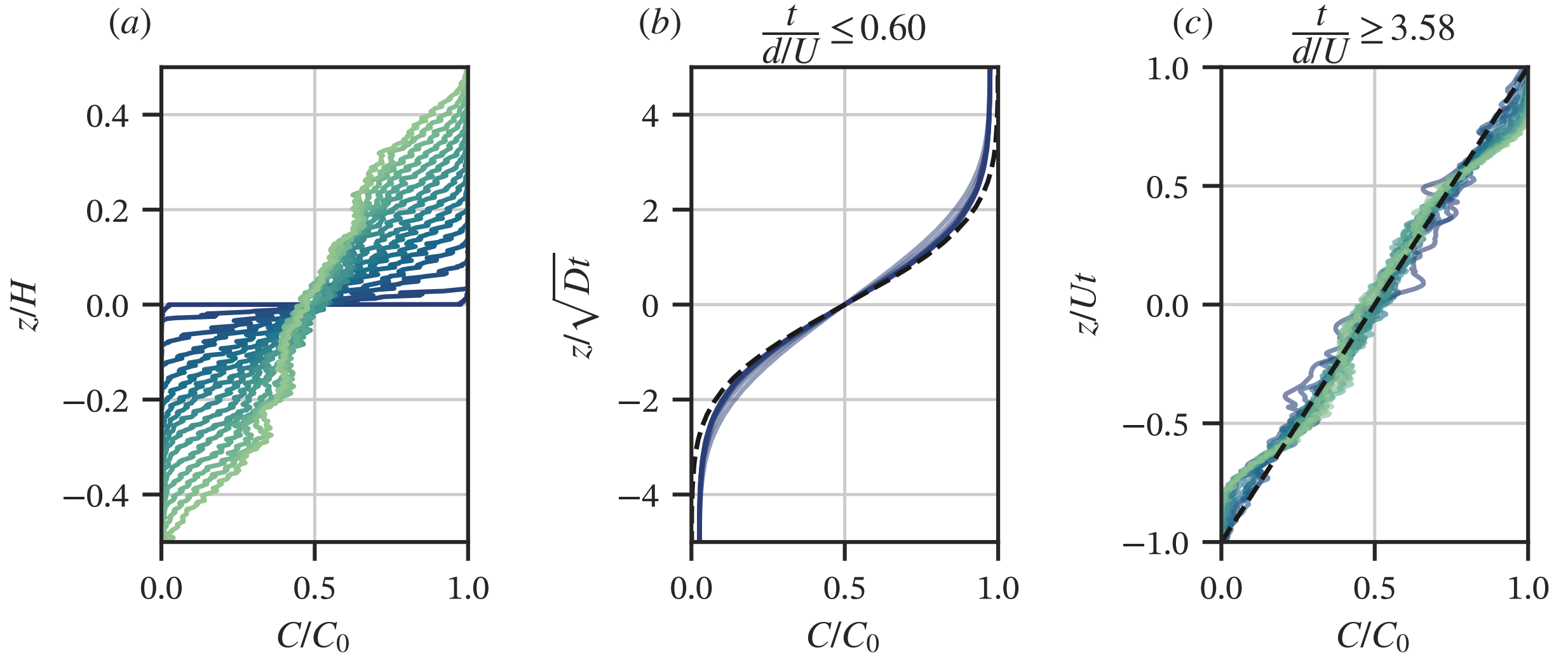
$$\begin{aligned} \partial_t \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u} &= -\rho_0^{-1} \nabla p + \nu \nabla^2 \mathbf{u} - g\beta C \hat{\mathbf{z}}, \\ \partial_t C + (\mathbf{u} \cdot \nabla) C &= D \nabla^2 C, \\ \rho &= \rho_0 \left[1 + \frac{\Delta \rho}{\rho_0 C_0} (C - C_0) \right] \end{aligned}$$



$$U = \frac{g\Delta\rho k}{\mu}$$

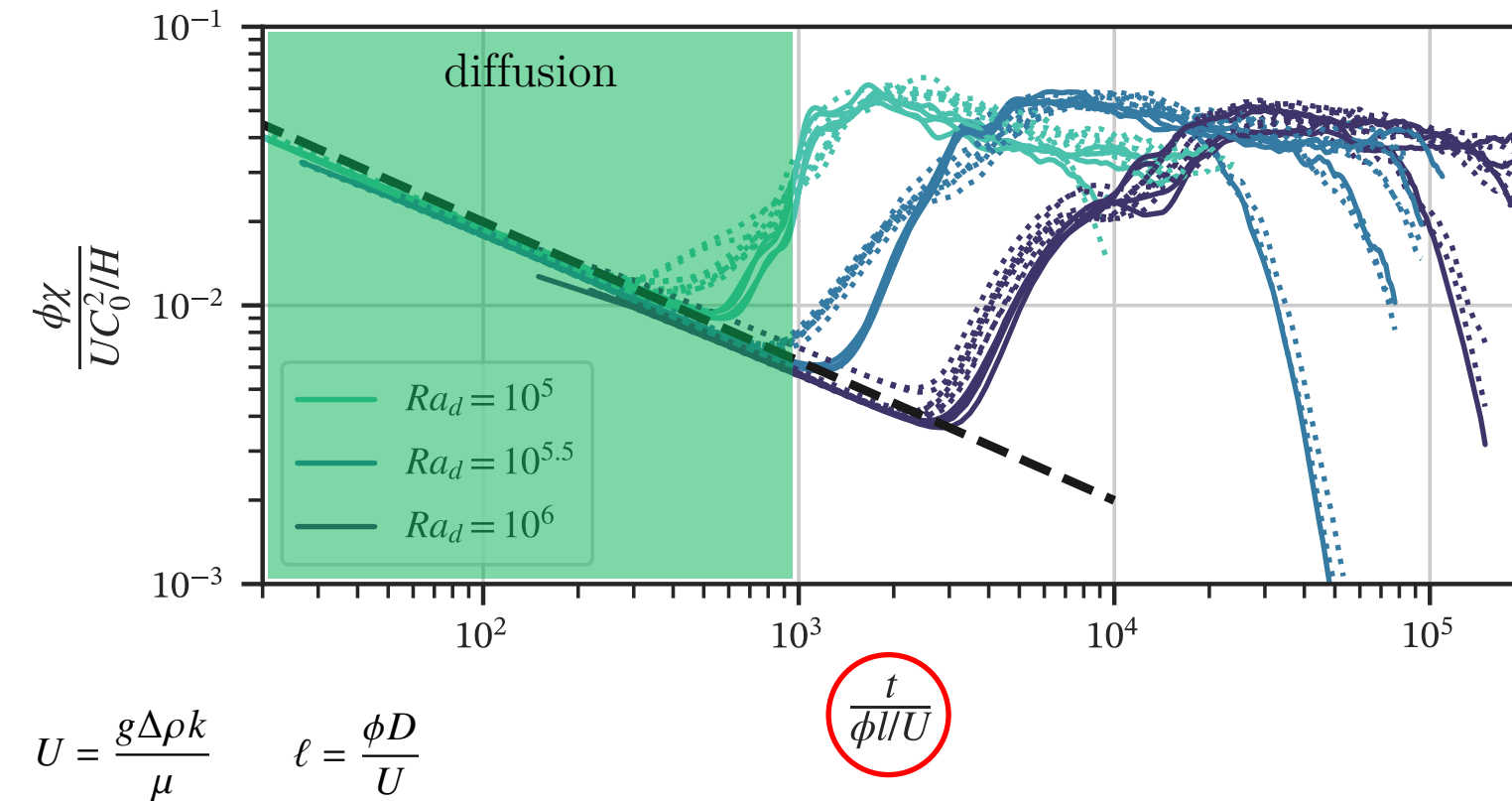
$$\ell = \frac{\phi D}{U}$$





$$\chi = D \langle |\nabla C|^2 \rangle_f = \frac{D}{V_f} \int_{V_f} |\nabla C|^2 dV$$

Can we model this mixing/dissolution process?



Diffusion:

$$C = C_0 + \frac{\Delta C}{2} \operatorname{erf} \left(\frac{z}{\sqrt{2\kappa t}} \right)$$

$$\partial_z C = \frac{\Delta C}{2\sqrt{\pi\kappa t}} \exp \left(-\frac{z^2}{2\kappa t} \right)$$

$$\chi = \kappa \langle |\nabla C|^2 \rangle = \frac{\kappa}{H} \int_{-\infty}^{\infty} |\partial_z C|^2 dz$$

$$= \sqrt{\frac{\kappa}{8\pi t}} \frac{(\Delta C)^2}{H}$$

$$\chi = D \langle |\nabla C|^2 \rangle_f = \frac{D}{V_f} \int_{V_f} |\nabla C|^2 dV$$

Convection

$$\chi = \kappa \langle |\nabla C|^2 \rangle = \kappa \frac{L_m}{H} \langle |\nabla C|^2 \rangle_{ML},$$

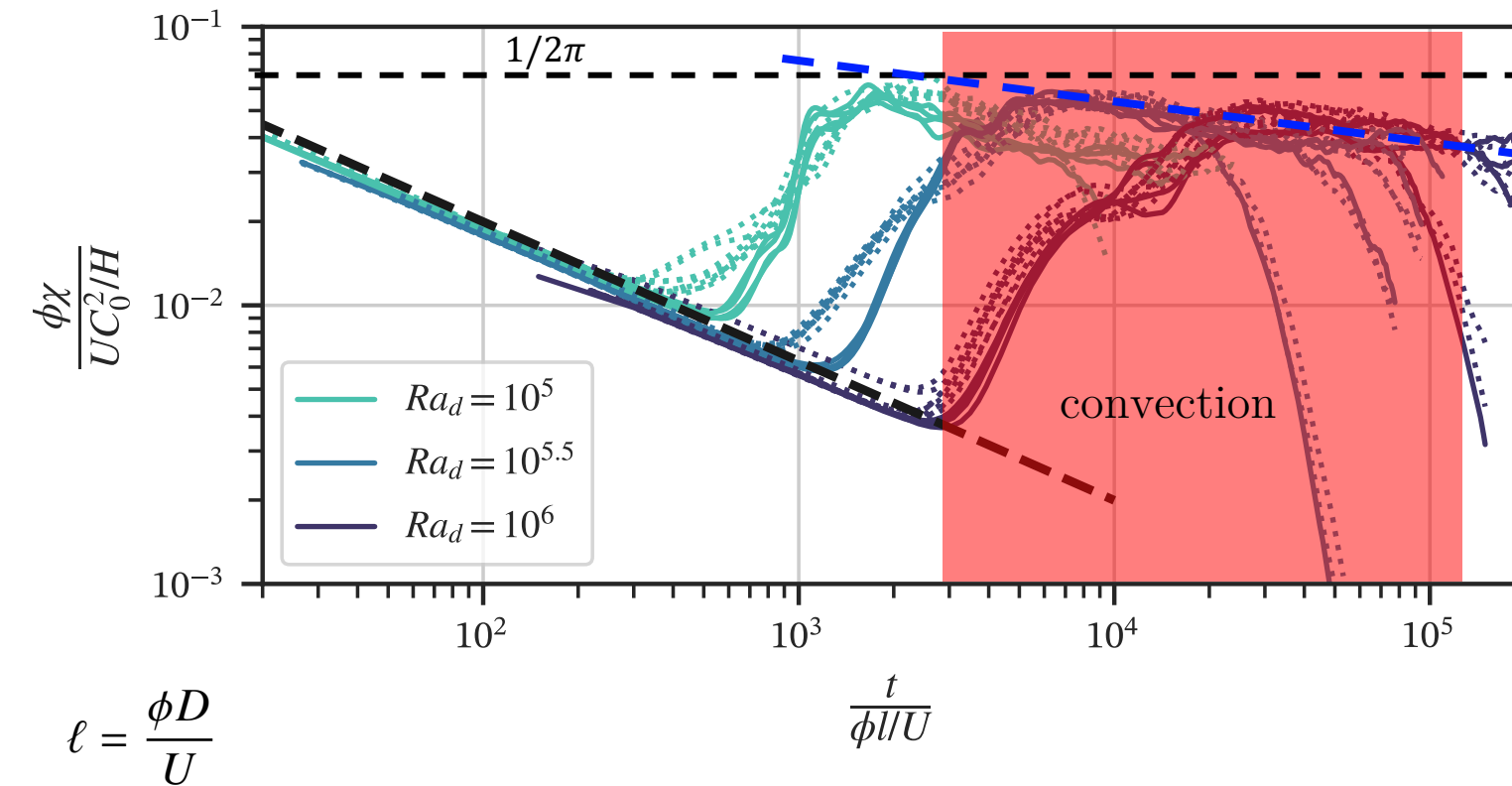
$$|\nabla C| \approx \frac{\Delta C}{2\sqrt{\pi \kappa t}}.$$

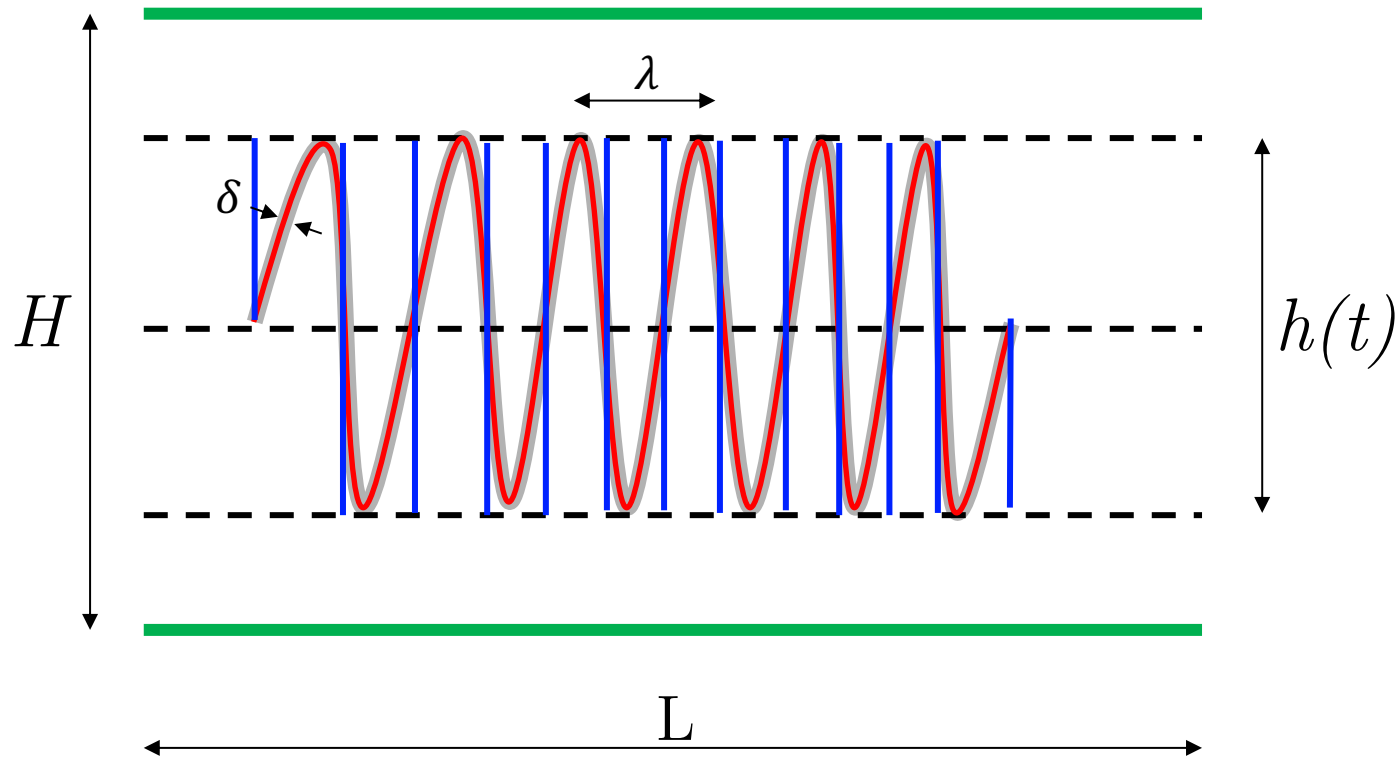
$$L_m \approx 2Ut,$$

$$\chi \approx \kappa \frac{2Ut}{H} \frac{(\Delta C)^2}{4\pi \kappa t} = \frac{1}{2\pi} \frac{U_d (\Delta C)^2}{H}.$$

$$\frac{\phi \chi}{UC_0^2/H} = 1/(2\pi)$$

$1/(2\pi)$ is the maximum value of mean dissipation. Measurements indicate that χ decreases with time





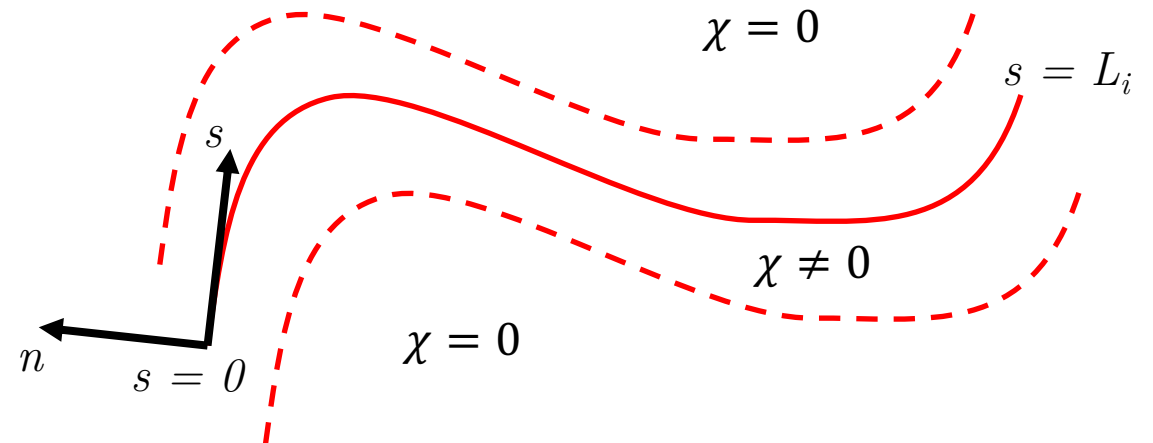
$$\chi = D \langle |\nabla C|^2 \rangle = \frac{D L_i}{H L} \int_{-\delta/2}^{+\delta/2} |\partial_n C|^2 dn$$

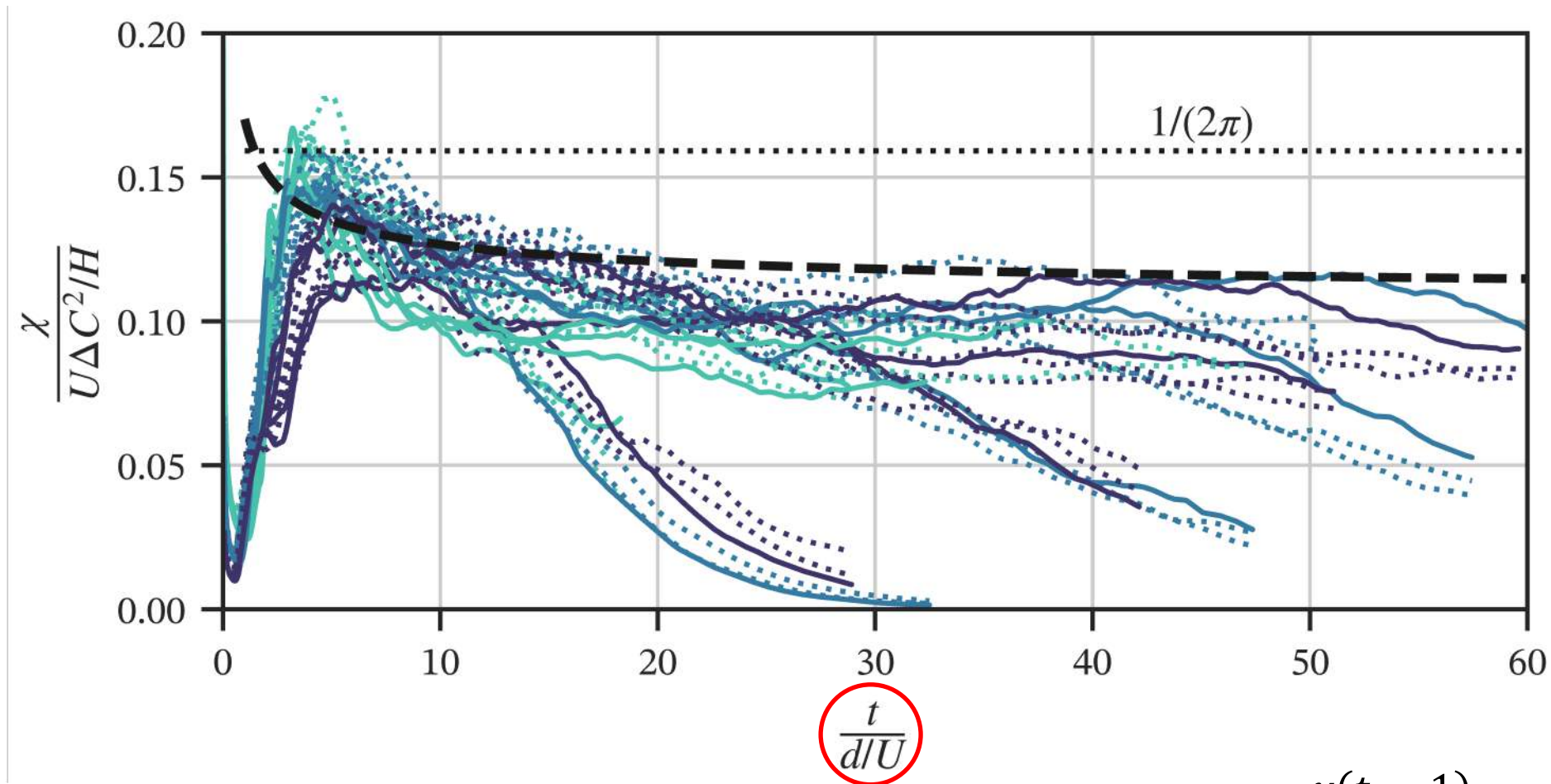
Assume:

1) Interface grows as:

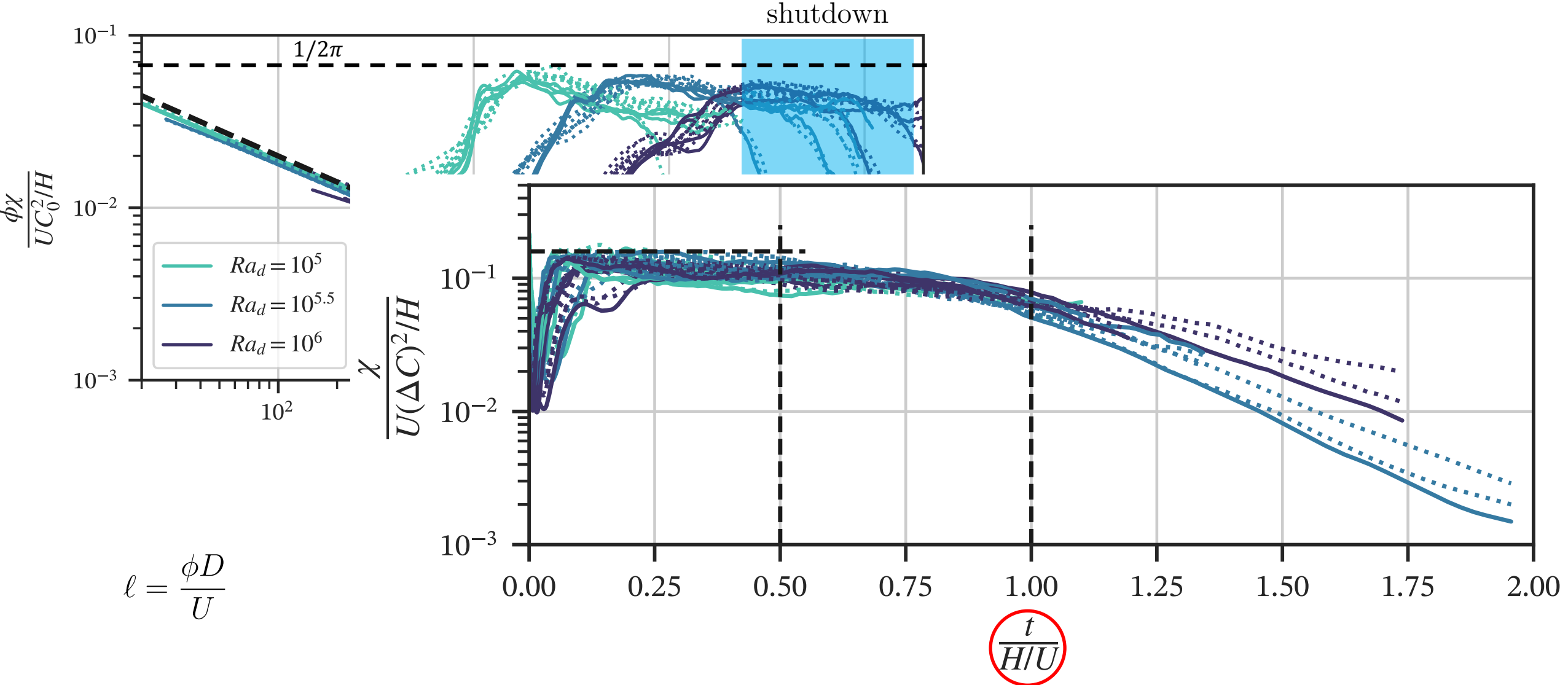
$$L_i = L + 2 N_{finger} h = L + 2 \frac{L}{\lambda} h$$

2) Gradient across the interface evolves according to the diffusive solution





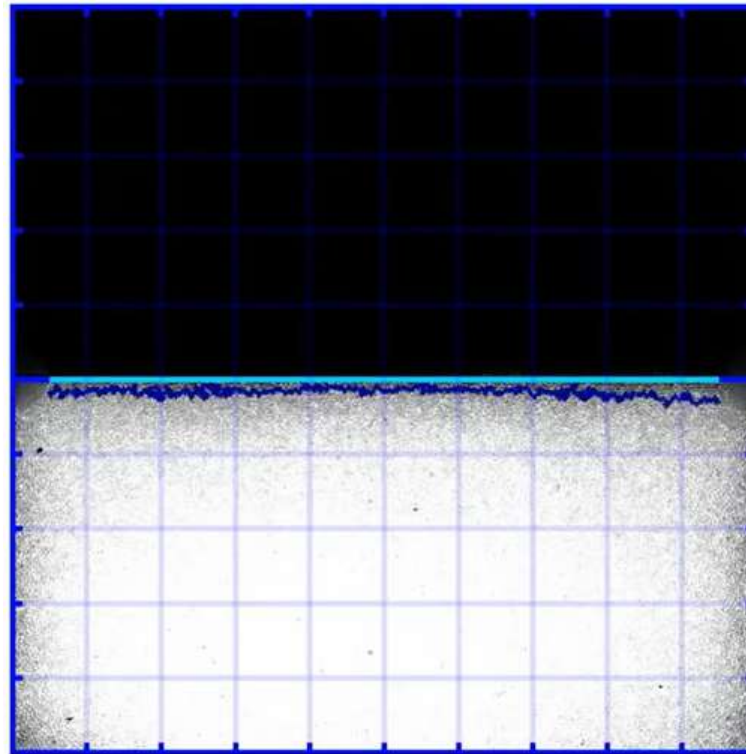
$$\frac{\chi(t=1)}{(\Delta C)^2 U / H} = \frac{\beta}{\alpha \pi} \left(1 + \frac{\alpha}{4} \right) \approx \frac{1}{1.92\pi} \approx \frac{1}{2\pi}$$



- Multiple **length scales** are relevant to **different phases** of the process
- We explain theoretically the scaling laws observed
- We plan to performed simulations in three-dimensional domains and **Darcy simulations** with **dispersion**



Der Wissenschaftsfonds.



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SURF

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High-resolution images, movies and slides are available upon request to m.depaoli@utwente.nl